

Quantum Mechanics

Quantisation $L = I\omega = m_e v_n r_n = n\hbar$

Balmer Series $\nu = \frac{m_e e^4}{8\epsilon_0^2 \hbar^3} \left(\frac{1}{2^2} - \frac{1}{m^2} \right)$

Heisenberg $\Delta x \Delta p \geq \frac{\hbar}{2}$

Dispersion $\omega = \frac{\hbar k^2}{2m}, \omega = ck$

Schrödinger $-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi = E\psi$
 $-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2} + V(x,t)\psi(x,t) = i\hbar \frac{\partial \psi(x,t)}{\partial t}$

Tunnelling Alpha decay (atom potential)
 Ammonia molecule (nitrogen atom penetrates plane of H)
 Tunnel diode (variable height voltage potential barrier)
 STM (tunnelling = distance)

Quantum Rotator $J = j(j+1)\hbar^2$

Quantum Oscillator $E_n = (n + \frac{1}{2})\hbar\omega$

Statistical Mechanics

Perfect Gas $pV = NkT, p = nkT$

First Law $dU = dQ + dW$

Constant V $dU = dQ = \left(\frac{\partial Q}{\partial T}\right)_V dT = C_V dT$

Constant P $\left(\frac{\partial Q}{\partial T}\right)_P = C_P = C_V + R$

Adiabatic $dQ = 0, pV^{\frac{C_P}{C_V}} = C$

Boltzmann $p(E)dE = Ag(E) \exp\left(-\frac{E}{kT}\right)dE$

Kinetic Theory Perfect gas (equation of state, no heat in a Joule expansion)
 Gases consist of small, hard, independent particles
 Velocity randomly distributed
 Temperature related to KE

1D Maxwell $f_1(v_x)dv_x = \sqrt{\frac{m}{2\pi kT}} e^{-mv_x^2/2kT} dv_x$

3D Maxwell $f(v)dv = \left(\frac{m}{2\pi kT}\right)^{\frac{3}{2}} 4\pi v^2 e^{-mv^2/2kT} dv$

$$\frac{1}{2}mv^2 = \frac{3}{2}kT, \bar{v} = \sqrt{\frac{8kT}{\pi m}}$$

Pressure $p = nm\bar{v}_x^2 = \frac{1}{3}nm\bar{v}^2$

Partial Pressures $\frac{\sigma_{M2}}{\sigma_{M1}} = \left(\frac{M_1}{M_2}\right)^{1/2}, p = p_1 + p_2$

Flux Density $J = \frac{\Delta N}{\Delta t} = n \int_0^\infty v_x f_1(v_x) dv_x = \frac{1}{4}n\bar{v}$

Mean Free Path $\lambda = \frac{1}{n\sigma}, p(x) = \exp\left(-\frac{x}{\lambda}\right), R = \lambda\sqrt{N}$

Equipartition $E = \frac{f}{2}kT, \left(\frac{\partial Q}{\partial T}\right)_V = C_V = \frac{f}{2}R$

Second Law No process exists whose sole effect is the transfer of heat from a cooler to a hotter body
 $\ln g = N \ln(Q/N), Q$ quanta, N boxes, g = permutations
 Net energy exchange leads to an increase in the most probable

number of ways in which the system can arrange itself

Entropy $S = k \ln g, \frac{1}{T} = \frac{\partial S}{\partial E}$

Quantum Statistics

Raleigh-Jeans Body Radiation $dN = \frac{8\pi\nu^2}{c^3} d\nu, u(\nu)d\nu = \bar{E}dN = \frac{8\pi\nu^2 kT}{c^3} d\nu$

Planck Black Body Radiation $u(\nu)d\nu = \frac{8\pi h\nu^3}{c^3} \frac{1}{\exp(h\nu/kT)-1} d\nu$

$$u = \int_0^\infty u(\nu)d\nu = \left(\frac{8\pi^5 k^4}{15c^3 h^3}\right) T^4$$

Schottky Anomaly A maximum in the heat capacity, where you run out of energy levels in which to store heat

Excitation Condition $T \gg \frac{E}{k}$

Einstein Solid Heat Capacity $C = 3R \left(\frac{h\nu_E}{kT}\right)^2 \frac{\exp(h\nu_E/kT)}{(\exp(h\nu_E/kT)-1)^2}$

Debye Solid Heat Capacity $C \propto VT^3$ at low T by finding photon energies that will fit into a solid